# ON THE STRONG STABLITY OF RESONANT SYSTEMS UNDER PARAMETRIC PERTURBATIONS 

PMM Vol.41, № 2, 1977, pp. 251-261<br>Ia.M. GOL'TSER<br>(Alma-Ata)<br>(Received August 14, 1976)

An autonomous system of differential equations with holomorphic right-hand sides that continuously depend on a parameter is considered. It is assumed that in the considered region of parameter values the first approximation system has $n$ pairs of pure imaginary roots.

The stability properties of a system may undergo various changes when parameter $\mu$ is varied near its resonance value $\mu_{0}$ at which the system has an internal resonance. The problem of strong stability is posed in the case when the character of the system stability at point $\mu_{0}$ remains unchanged in some neighborhood of that point.

A normal form, continuous with respect to $\mu$ of the system is derived. The form, unlike the usual and normal form [1,2], does not change its structure at transition of the system through resonance (*). Conditions of strong asymptotic stability and instability are obtained for resonance of odd order. Cases of "explosive instability" in which the asymptotic stability in the resonance point neighborhood changes into instability at point $\mu_{0}$ are separated out.

The concept of the present work is akin to that of [3], where the problem of stability in the presence of parametric perturbations is stated and substantiated. Several aspects of this problem and its generalization were discussed in [4]. The allied question of the form of the problem of the difference between dangerous and safe limits of the stability region was considered in [5].

1. Statement of the probiem. Let us consider the system of differential equations

$$
\begin{equation*}
z^{\cdot}=P(\mu) z+\sum_{l=k \geqslant 2}^{\infty} Z^{(l)}(z, \mu), \quad \mu \in\left(\mu_{1}, \mu_{2}\right)=D \tag{1.1}
\end{equation*}
$$

where $z$ is an $r$-dimensional vector, $P(\mu)$ is an $(r \times r)$-matrix, $Z^{(l)}(z, \mu)$ are vector-forms of $l$-th order with respect to $z$. Matrix $P(\mu)$ and the coefficients of form $Z^{\prime \prime}(z, \mu)$ are continuous functions of parameter $\mu$.

We shall call system (1.1) stable (unstable) at point $\mu_{0} \in D$, if the zero solution of that system for $\mu=\mu_{0}$ is Liapunov stable (unstable).

The considered problem is associated with the analysis of the effect of small variation of the parameter on the stability properties of system (1.1).

Definition 1. 1. System (1.1) which is stable (unstable) at point $\mu_{0} \in D$ is called strongly stable (unstable) at that point, if there exists such $\varepsilon$-neighborhood

[^0]$U_{\mathrm{E}}\left(\mu_{0}\right) \subset D$ of point $\mu_{0}$ that system (1.1) is stable (unstable) at every point $\mu \in$ $U_{e}\left(\mu_{0}\right)$.

Note 1.1. The term "strong stability" conforms to the terminology used in investigations of linear periodic Hamiltonian systems for small (but more general than parametric) perturbations of the Hamiltonian [ [i].
System (1.1) is considered below on the assumption that matrix $P(\mu)$ has in region I) $n$ pairs of continuous with respect to $\mu$ pure imaginary eigenvalues, and is reducible in $D$ to a diagonal matrix which is continuous with respect to $\mu$ (The question of the normal form of matrices that depend on a parameter was investigated in [7]).

We set $r=2 n$ and denote the eigenvalues of matrix $P(\mu)$ by $\pm \lambda_{s}(\mu), s=1$, $2, \ldots, n, \quad \lambda_{s}{ }^{2}(\mu)<0(\forall \mu \in D)$, and by $R_{n}{ }^{+}$the set of $n$-dimensional vectors whose components are nonnegative integers. If $m \in R_{n}{ }^{+}$, then $m=\left(m_{1}, \ldots, m_{n}\right)$, $m_{s} \in R_{1}^{+}$, and $|m|=m_{1}+\ldots$ i- $m_{n}$.

Definition 1.2. At point $\mu_{0} \in D$ system (1.1) has a $j$-th order internal resonance, if there exists vector $m \in R_{n}{ }^{+}$with $m_{\text {s }}$ relatively prime, such that when $\mu=$ $\mu_{0}$

$$
\begin{equation*}
\left\langle m, \lambda^{\circ}\right\rangle=\sum_{s=1}^{n} m_{s} \lambda_{s}^{\circ}=0, \quad|m|=j, \quad \lambda^{\circ}=\lambda\left(\mu_{0}\right)=\left(\lambda_{1}^{\circ}, \ldots, \lambda_{n}^{\circ}\right) \tag{1.2}
\end{equation*}
$$

An effective method of investigating the stability of system (1.1) with fixed values of the parameter (either resonant or nonresonant) consists of its preliminary reduction to the normal form up to terms of a reasonably high order [1,2]. If point. $\mu_{0}$ is nonresonant, such normalization in a fairly small neighborhood of it is continuous with resect to $\mu$, and one can expect that the system stability properties are retained in that neighborhood. When parameter $\mu$ passes through its resonance value $\mu_{0}$, the structure of the usual normal form considerably changes. Hence for the considered class of systems the problem of strong stability in the neighborhood of resonant values of the parameter is of considerable interest. The question of existence for system (1.1) of a normal form continuous with respect to $\mu$, which arises in this connection, is considered below.
2. The continuous normal form. On the assumptions made in Sect, 1 , system (1.1) may be written as (the upper dash denotes a complex-conjugate quantity)

$$
\begin{align*}
& x^{\cdot}=\Lambda(\mu) x+\sum_{l=k \geqslant 2}^{\infty} X^{(l)}(x, y, \mu)  \tag{2.1}\\
& y^{\cdot}=-\Lambda(\mu) y+\sum_{l=k \geqslant 2}^{\infty} Y^{(l)}(x, y, \mu)
\end{align*}
$$

where $x=\bar{y}=\left(x_{1}, \ldots, x_{n}\right), E$ is a unit $(n \times n)$-matrix, and $\Lambda(\mu)=\lambda(\mu) E$. We represent the components of $n$-dimensional vector-forms $X^{(l)}=Y^{(l)}$ in the form

$$
X_{3}^{(l)}(x, y, \mu)=\sum_{|p|+|q|=} a_{p, q}^{s}(\mu) x^{p} y^{q}, \quad p, q \in R_{n}^{+}, x^{p}=x_{1}^{p_{1}} \ldots x_{n}^{p_{n}}
$$

Functions $\lambda_{s}(\mu)$ and $a_{p, q}^{\prime}(\mu)$ are continuous in $D$.
Let us consider system

$$
\begin{equation*}
u^{\cdot}=\Lambda(\mu) u+\sum_{l=k \geqslant 2}^{\infty} U^{(l)}(u, v, \mu) \tag{2.2}
\end{equation*}
$$

$$
\begin{aligned}
& v^{\cdot}=-\Lambda(\mu) v+\sum_{l=k \geqslant 2}^{\infty} V^{(l)}(u, v, \mu) \\
& U_{s}^{(l)}=\bar{V}_{s}^{(l)}=\sum_{|p|+|q|=i} \alpha_{p, q}^{s}(\mu) u^{p} v^{q}
\end{aligned}
$$

and transformation

$$
\begin{align*}
& x=u+\sum_{l=k \geqslant 2}^{\infty} \Phi^{(l,}(u, v, \mu), \quad y=v+\sum_{l=k \geqslant 2}^{\infty} \bar{\Phi}^{(l)}(u, v, \mu)  \tag{2.3}\\
& \Phi_{s}^{(l)}=\sum_{|p|+|q|=l} A_{p, q}^{s}(\mu) u^{p} v^{q}
\end{align*}
$$

together with (2.1), and seek the simplest form of system (2.2) to which system (2.1) can be reduced by the transformation, continuous with respect to $\mu$, in the form of formal series (2.3).

For the successive determination of forms $\Phi_{s}{ }^{(l)}$ from (2.1) - (2.3) we obtain

$$
\begin{align*}
& \sum_{r=1}^{n} \lambda_{r}\left(\frac{\partial \Phi_{s}^{(l)}}{\partial u_{r}} u_{r}-\frac{\partial \Phi_{s}^{(l)}}{\partial v_{r}} v_{r}\right)=\lambda_{s} \Phi_{s}^{(l)}-U_{s}^{(l)}-F_{s}^{(l)}+X_{s *}^{(l)}  \tag{2,4}\\
& F_{s}^{(l)}=\sum_{r=1}^{n} \sum_{j=k}^{l-1} \frac{\partial \Phi_{s}^{(j)}}{\partial u_{r}} U_{r}^{(l-j+1)}+\frac{\partial \Phi_{s}^{(j)}}{\partial v_{r}} V_{r}^{(l-j+1)}
\end{align*}
$$

where $X_{s}{ }^{(l)}$ is an $l$-th order form in the expansion in series of functions

$$
\sum_{j=k \geqslant 2}^{i} X_{s}^{(j)}(u+\ldots, v+\ldots)
$$

For the simultaneous determination of coefficients of forms $\Phi_{s}{ }^{(l)}$ and $U_{s}{ }^{(l)}$ we obtain

$$
\left\langle p-q-\delta_{s}, \lambda(\mu)\right\rangle A_{p, q}^{s}=a_{p, q}^{s *}(\mu)-f_{p, q}^{s}(\mu)-\alpha_{p, q}^{s}(\mu)
$$

where.$a_{p, q}^{s *}(\mu)$ and $f_{p, q}^{s}(\mu)_{t}(|p|+|q|=l)$ are coefficients of forms $X_{s *}^{(l)}$ and $F_{s}{ }^{(l)}$, and $\delta_{s}=(0, \ldots, 1, \ldots, 0)$ is the basis vector in $R_{n}{ }^{+}$.

Definition 2.1. The bivectors $(p, q), p \neq q+\delta_{s}$ and the related coefficients $s$ - $x$ in the equations of systems (2.1) and (2.2), and in the transformation (2.3) are called resonant, if there exists a $\mu_{0} \in D$ such that

$$
\left\langle p-q-\delta_{s}, \lambda\left(\mu_{0}\right)\right\rangle=0
$$

We denote the set of all resonant bivectors of the $s$-th equation by $L_{D^{s}}$. When $p=$ $q+\delta_{\text {s }}$ for the bivectors ( $p, q$ ) we have

$$
\left\langle p-q-\delta_{s}^{\prime}, \lambda(\mu)\right\rangle \equiv 0 \quad(\forall \mu \in D)
$$

which constitute the set $L_{0}{ }^{s}$ that determines the terms of the identical resonance. We denote $L^{s}=L_{D}{ }^{s} \cup L_{0}{ }^{s}$.

Let functions $a_{p, q}^{\mathrm{s}}$ and $f_{p, q}^{\mathrm{s}}$ in (2.5) be continuous in $D$. It is evident that when ( $p$, $q) \in L^{s}$, then (2.5) has a solution that is continuous with respect to $\mu$ for any selected continuous functions $\alpha_{p, q}^{s}(\mu)$.

If $(p, q) \in L_{D}{ }^{\prime}$, there exists a point $\mu_{0} \in D$ such that $\left\langle p-q-\delta_{s}, \lambda\right\rangle \rightarrow 0$ when $\mu \rightarrow \mu_{0}$. For function $A_{p, q}^{s}(\mu)$ to be continuous in $D$ it is necessary to determine function $\alpha_{p, q}^{s}(\mu)$ so that it is continuous in $D$, and that at each resonant point
of region $D$ the right-hand side of (2.5) is infinitely small of an order not less than $\left\langle p-q-\delta_{s}, \lambda(\mu)\right\rangle$. It is obvious that with such choice the structure of functions $\alpha_{p, q}^{s}$ and $A_{p . q}^{s}$ in the neighborhood of the resonance point is of the form

$$
\begin{align*}
& \alpha_{p, q}^{s}(\mu)=a_{p, q}^{s *}\left(\mu_{0}\right)+f_{p, q}^{s}\left(\mu_{0}\right)-\varepsilon_{p, q}^{s}(\mu)  \tag{2.6}\\
& A_{p, q}^{s}(\mu)=\left\langle p-q-\delta_{s}, \quad \lambda(\mu)\right\rangle^{-1} \varepsilon_{p, q}^{s}(\mu), \quad(p, \quad q) \in L_{D}^{s}
\end{align*}
$$

where $\varepsilon_{p, q}^{s}$ is an arbitrary continuous function in $D$ such that there exists the limit

$$
\begin{equation*}
\lim _{\mu \rightarrow \mu_{0}} A_{p, q}^{s}(\mu)=A_{p, q}^{s}\left(\mu_{0}\right) \tag{2.7}
\end{equation*}
$$

The constant $A_{p, q}^{s}\left(\mu_{0}\right)$ can be arbitrarily chosen, and this can be used for further simplification of the system in the case of third order resonances.

If $(p, q) \in L_{0}{ }^{s}$, then it is sufficient to set in (2.5)

$$
\begin{equation*}
\alpha_{q+\delta_{\mathbf{s}}, q}^{s}=a_{q+\delta_{s}, q}^{s *}(\mu)+f_{p, q}^{s}(\mu) \tag{2.8}
\end{equation*}
$$

Taking now into consideration that forms $X_{s *}{ }^{(l)}$ and $F_{s}{ }^{(l)}$ depend only on forms $\Phi_{s}{ }^{(j)}$ and $U_{s}{ }^{(j)}$, where $j<l$, and that $X_{s *}{ }^{(l)}=X_{s}{ }^{(l)}$ and $F_{s}{ }^{(l)}=0$ when $l=k$, it is possible to assert that functions $f_{p, q}^{s}$ and $a_{p, q}^{8}$ are continuous in $D$ for all $l=k, k+$ $1, \ldots$ Hence the following theorem is valid.

Theorem 2.1. With the arbitrary selection in system (2.2) that are continuous with respect to $\mu$ it is possible to select for that system continuous resonant coefficients so that the continuous in $D$ transformation (2.3), which transforms system (2.1) into(2.2), is obtained.

The continuous normal form is obtained by equating in (2.2) all nonresonant coefficients to zero. The structure of the continuous normal form is then

$$
\begin{align*}
& u_{\mathrm{s}}^{*}=\lambda_{s}(\mu)+u_{s} \sum_{q=[k / 2] \geqslant 1}^{\infty} \alpha_{q+\delta_{s}, q}^{s} \omega^{q}+\sum_{(p, q) \in L_{D}^{s}} \alpha_{p, q}^{s} u^{p} v^{q}  \tag{2.9}\\
& v_{s}^{*}=\bar{u}_{\mathrm{s}}^{*}, \quad \omega=\left(\omega_{1}, \ldots, \omega_{n}\right), \quad \omega_{s}=u_{\mathrm{s}} v_{\mathrm{s}}
\end{align*}
$$

The following properties of the transformation and of the normal form, which will be used subsequently, should be noted. If $k \geqslant 2$ is the lowest order of nonlinear terms in (2.1), forms $\Phi_{s}{ }^{(j)}$ and $U_{s}{ }^{(j)}$ for $k \leqslant j \leqslant 2 k-2$ are determined independently of forms $\Phi_{s}{ }^{(l)}$ and $U_{s}{ }^{(b)}$ for $l<j$. This property holds in conventional normalization and is due to the structure of functions $X_{s *}{ }^{(l)}$ and $F_{s}{ }^{(l)}$ in (2.4).

Let us compare the structure of system (2.9) for some fixed $\mu_{0}$ with that of the usual normal form at that point. We use the notation $L_{\mu_{0}}{ }^{s}=\left\{(p, q) \mid\left\langle p-q-\delta_{s}, \lambda^{\circ}\right\rangle=0\right\}$. It is clear that $L_{\mu_{0}}{ }^{\text {s }} \subset L_{D}{ }^{8}$.

In the case of conventional normalization of system (2,1) the $s$-th equation contains only terms corresponding to bivectors $(p, q) \in L_{\mu_{0}}{ }^{s} \cup L_{0}{ }^{s}$ when $\mu=\mu_{0}$, while system (2.9) contains in addition to these terms also those for which $(p, q) \in L_{D}{ }^{s} \backslash L_{\mu_{0}}{ }^{s}$.

The two forms coincide at point $\mu_{0}$ only when $L_{D}{ }^{s} \backslash L_{\mu_{0}}{ }^{s}=\varnothing$, but in the neighborhood of point $\mu_{0}$, where $\mu \neq \mu_{0}$, they are different. Conventional normalization is continuous only when $L_{D}{ }^{s}=\varnothing$ (as noted in Sect. 1).

When solving the problem of strong stability at point $\mu_{0}$ we consider a reasonably small $\varepsilon$-neighborhood of point $\mu_{0}$ as the region $D=U_{\varepsilon}\left(\mu_{0}\right)$. We denote the deleted neighborhood $U_{\varepsilon}\left(\mu_{0}\right)$ of point $\mu_{0}$ by $D^{*}=U_{\varepsilon}^{*}$.

Let point $\mu_{0}$ be resonant whose unique lowest resonance is of order $g$. We assume (and this is the general case) that $\mu_{0}$ is the isolated root of the equation

$$
\langle m, \lambda(\mu)\rangle=0 \text { for }|m|=g
$$

The number $\varepsilon$ is assumed to be so small, that when this equation has in $D$ other integral solutions, their norm is $|m| \gg g$.

Under these conditions the lowest resonating terms in $D$ are of order $g-1$, hence the structures of the conventional and the continuous normal forms coincide (for $k<$ $g-1)$ to within terms of order $g-2$. If the solution of the problem of stability in region $D$ is independent of terms of order higher than $g-2$, the presence of inner resonance at point $\mu_{0}$ does not affect the solution of the problem of strong stability. To detect possible bifurcations of the properties of stability by the presence in the system of a resonance of order $g$, we assume that in (2.1) $k=g-1$, and that $k=2 N$.

On these assumptions, and restricting the reduction of the continuous normal form to terms of order up to and including the $(k+1)$-st, we obtain the following system:

$$
\begin{gather*}
u_{s}^{\cdot}=\lambda_{s}(\mu) u_{s}+\alpha_{s}(\mu) v^{m-\delta_{s}}+u_{s} \sum_{|p|=N} \alpha_{s}^{p}(\mu) \omega^{p}+O_{\mu}\left(\|u+v\|^{N+2}\right)  \tag{2.10}\\
v_{s}^{*}=\bar{u}_{s}^{\cdot}, \quad m, p \in R_{n}^{+}, \quad|m|=k+1=2 N+1 \geqslant 3
\end{gather*}
$$

where $\lambda_{s}(\mu), \alpha_{s}(\mu), \alpha_{s}{ }^{p}(\mu)$ and $O_{\mu}$ are continuous in $D$.
For $\mu=\mu_{0}$ system (2.10) is the same as that analyzed in [2,8] and is obtained from (2.9) by considering the structure of lower resonating terms similar to those in [2].
3. Some properties of the model iystem. Let us consider the model system in the $k-$ th approximation

$$
\begin{equation*}
u_{\mathrm{s}}^{*}=\lambda_{\mathrm{s}}(\mu) u_{s}+\alpha_{\mathrm{s}}(\mu) v^{m-\delta_{s}}, \quad v_{s}^{*}=-\lambda_{s}(\mu) v_{s}+\bar{\alpha}_{s}(\mu) u^{m-o_{s}} \tag{3.1}
\end{equation*}
$$

We introduce ancilliary angles $\varphi_{s}(\mu)$ setting

$$
\begin{equation*}
\sin \varphi_{s}=-a_{s}\left|\alpha_{s}\right|^{-1}, \quad \cos \varphi_{s}=b_{s}\left|\alpha_{s}\right|^{-1}, \quad \alpha_{s}=a_{s}+i b_{\mathrm{s}} \tag{3.2}
\end{equation*}
$$

and identify these with points of the unit trigonometric circle. Obviously $\varphi_{s}(\mu)$ are continuous functions of $\mu$ in $D$.

Let $d_{s}$ be the diameter of the circle drawn through point $\varphi_{s}$. Two incompatible dispositions of points $\varphi_{s}$ are possible.
A. For any diameter $d_{s}$ there exist points $\varphi_{j}$ that lie on different sides of $d_{s}$.
B. There exists diameter $d_{s}$ such that all $\varphi_{j}$ points lie on one side of it.

Case A necessitates that $n \geqslant 3$, while case B always obtains when $n=2$. If case A obtains at point $\mu_{0}$ (of region $D$ ), we say that conditions $A(\mu)(A(D))$ are satisfied. Similarly $B(\mu), B(D)$.

The algebraic characteristic of case $A$ is defined by the following lemma.
Lemma 3.1. For condition $A(\mu)$ to be satisfied it is necessary and sufficient if there exist such coefficients $\alpha_{p_{1}}(\mu), \alpha_{s_{2}}(\mu)$ and $\alpha_{s_{1}}(\mu)$ that

$$
\operatorname{sign} D_{s_{152}}=\operatorname{sign} D_{s_{2} s_{3}}=-\operatorname{sign} D_{s_{1} s_{3},} \quad D_{s_{j} s_{k}}=\left|\begin{array}{ll}
a_{s_{j}} & a_{s_{k}}  \tag{3.3}\\
b_{s_{j}} & b_{s_{k}}
\end{array}\right|
$$

The following lemma is also valid.

Lemma 3.2. For condition $B(\mu)$ to be satisfied it necessary and sufficient that there exists such numbering of angles $\varphi_{s}(\mu)$ for which

$$
\begin{equation*}
\varphi_{1}(\mu) \leqslant \varphi_{2}(\mu) \leqslant \cdots \leqslant \varphi_{n}(\mu) \leqslant \varphi_{1}(\mu)+\pi \tag{3.4}
\end{equation*}
$$

With the use of inequalities (3.4) we subdivide case $B$ into the following subcases:

$$
\begin{aligned}
& \left.B_{1} \quad \text { (鸟 } l \mid 1 \leqslant l \leqslant n-1\right) ~\left(\varphi_{1}=\ldots=\varphi_{l}<\varphi_{l+1} \leqslant \ldots .\right. \\
& \left.. \leqslant \varphi_{n}<\varphi_{1}+\pi\right) \\
& B_{2}(\forall s)\left(\varphi_{s}=\varphi_{1}\right) \\
& B_{3}\left(\left.\begin{array}{l}
0
\end{array} \right\rvert\, 1 \leqslant k \leqslant n-1\right) \quad\left(\varphi_{1}=\ldots=\varphi_{k}<\varphi_{k+1}=\ldots\right. \\
& \left.=\varphi_{n}=\varphi_{1}+\pi\right) \\
& B_{4} \quad\left(\Psi_{j} \mid 3 \leqslant j \leqslant n\right) \quad\left(\varphi_{1} \leqslant \varphi_{2} \leqslant \cdots<\varphi_{j-1}<\varphi_{j}=\ldots\right. \\
& \left.=\varphi_{n}=\varphi_{1}+\pi\right)
\end{aligned}
$$

It is possible to show that condition $A(\mu)$ is equivalent to the following statement of a geometrical nature.

Lemma 3.3. For condition $A(\mu)$ to be satisfied it is necessary and sufficient that there exist such points $\varphi_{s 1}(\mu), \varphi_{s 2}(\mu)$ and $\varphi_{s 1}(\mu)$ that the triangle formed by these is acute.

Condition (3.3) is the algebraic criterion of acuteness of $\triangle \varphi_{s_{1}} \varphi_{s_{2}} \varphi_{s_{3}}$. In case $B_{1}$ all nondegenerate triangles are abtuse, and in case $B_{4}$ there is at least one right angle among them.

It follows from the above lemmas (without proof) and from the continuity of functions $\varphi_{s}(\mu)$ that when conditions $A\left(\mu_{0}\right)$ and $B_{1}\left(\mu_{0}\right)$ are satisfied, then for a fairly small $\varepsilon$ conditions $A(D)$ and $B_{1}(D)$ are also satisfied. Conditions $B_{2}\left(\mu_{0}\right), B_{3}\left(\mu_{0}\right)$ and $B_{4}\left(\mu_{0}\right)$ cannot be maintained when $\mu$ is varied. Case $B_{2}$ may convert to $B_{1}$, while $B_{3}$ and $B_{4}$ to $A$ or $B_{1}$.

As shown in $[2,8]$ system (3.1) has in the majority of cases when $\mu=\mu_{0}$, a set of integrals of the form

$$
\begin{equation*}
V_{0}=\sum_{s=1}^{n} \gamma_{s}{ }^{\circ} \omega_{s}, \quad \omega_{\mathrm{s}}=u_{s} v_{s}, \quad \gamma_{\mathrm{s}}{ }^{\circ}=\mathrm{const} \tag{3.5}
\end{equation*}
$$

The necessary and sufficient condition of stability of system (3.1) when $\mu=\mu_{0}$ is the presence among (3.5) of integrals of fixed sign. In the case of instability all integrals (3.5) are will alternating signs. In case $B_{1}$ when $n-2$ the system is unstable and has no integrals of the form (3.5). The necessary and sufficient condition of existence among (3.5) of fixed sign integrals is the fulfilment of condition $A\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$. In the unstable case $B_{4}$ there are fixed sign integrals among (3.5).

Let us consider matrix $C$ and the $n$-dimensional row vectors $a$ and $b$

$$
C=\left\|\begin{array}{c}
a_{1} a_{2}, \ldots a_{n} \\
b_{1} b_{2} \ldots b_{n}
\end{array}\right\|, \quad a=\left(a_{1}, \ldots, a_{n}\right), \quad b=\left(b_{1}, \ldots, b_{n}\right)
$$

The equality $D_{s_{j}{ }^{s}{ }_{k}}=\left|\alpha_{s_{j}} \alpha_{s_{k}}\right| \sin \left(\varphi_{s_{k}}-\varphi_{s_{j}}\right)$ implies that when conditions $A(\mu), B_{1}(\mu)$ and $B_{4}(\mu)$ are satisfied, we have Rank $C(\mu)=2$, while in cases $B_{2}$ and $B_{3}$ we have Rank $C(\mu)=1$.
Theorem 3.1. If matrix $C$ maintains its rank in region $D$ and the number of nonzero components of vector $a$ or $b$ exceeds the rank of that matrix, system (3.1) has the following set of integrals continuous with respect to $\mu$ :

$$
\begin{equation*}
V(\mu, \omega)=\sum_{s=1}^{n} \gamma_{s}(\mu) \omega_{3}, \quad \mu \in D, \quad V\left(\mu_{0}, \omega\right)=V_{0} \tag{3.6}
\end{equation*}
$$

The necessary and sufficient condition of the presence among these of integrals that are continuous in $D$ is the fulfilment of condition $A\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$.

Proof. First of all it is possible to ascertain that the assumption about the rank of the matrix is satisfied when conditions $A\left(\mu_{0}\right), B_{1}\left(\mu_{0}\right)$ and $B_{4}\left(\mu_{0}\right)$ are satisfied. It follows from this assumption that when conditions $B_{2}\left(\mu_{0}\right)$ and $B_{3}\left(\mu_{0}\right)$ are satisfied, then conditions $B_{2}(D)$ and $B_{3}(D)$ are satisfied, since the possible conversion of cases $B_{2}$ and $B_{3}$ to $A$ or $B_{1}$ for $\mu \neq \mu_{0}$ is associated with a change of the rank of matrix $C$.'

We determine the derivative of function $V$ by $(3.6)$ and by virtue of system (3.1) and from the equation $V^{*} \ldots 0$ we obtain for the determination of $\gamma_{s}(\mu)$ the system of equations

$$
\begin{equation*}
C^{\prime}(\mu) \gamma(\mu)=0, \quad \gamma(\mu)-\operatorname{colomn}\left|\gamma_{1}(\mu), \ldots, \gamma_{n}(\mu)\right| \tag{3.7}
\end{equation*}
$$

which is compatible with conditions of the theorem.
Let initially Rank $C(\mu)=2$. We denote vectors $a, b$ and $\gamma$ in which there are no components with subscripts $s_{1}, s_{2}$ and $s_{3}$ by $a^{\prime}, b^{\prime}$ and $\gamma^{\prime}$. The general solution of $(3,7)$ can be represented in the form

$$
\begin{align*}
& \gamma_{s_{1}}=D_{s_{2} b_{2}}^{-1}\left(\gamma_{s_{3}} D_{s_{4} s_{2}}+D_{s_{2} \rho}\right), \quad \gamma_{s_{2}}=D_{s_{s_{s}}}^{-1}\left(\gamma_{s_{s}} D_{s_{2} s_{2}}+D_{s_{2} \rho}\right)  \tag{3.8}\\
& \left(D_{s_{1} \div 2} \neq 0, \quad D_{s_{j} \rho}=\left|\begin{array}{cc}
a^{\prime} s_{j} & \rho_{1} \\
b_{s_{j}} & \rho_{2}
\end{array}\right|, \quad \rho_{1}=-a^{\prime} \gamma^{\prime}, \quad \rho_{2}=-b^{\prime} \gamma^{\prime}\right)
\end{align*}
$$

where $s_{1}$ and $s_{2}$ are selected so that the conditions appearing in parentheses are satisfied, and the components of vectors $\gamma^{\prime}$ and $\gamma_{s i}$ are free parameters of the solution,

If conditions $A\left(\mu_{0}\right)$ are satisfied, there exist such $s_{1}, s_{2}$ and $s_{3}$ for which formulas (3.3) valid at point $\mu_{0}$ are preserved also in region $D$. It follows directly from (3.8) that positive solutions of (3.7) exist only when conditions (3.3) are satisfied,

The free parameters which provide strictly positive solutions of the system (and by the same token fixed sign integrals of set (3.6)) must satisfy conditions

$$
\begin{equation*}
\gamma_{s}>0 \quad\left(s \neq s_{1}, s_{2}, s_{3}\right), \quad \gamma_{s_{3}}>\max \left\{D_{s_{1} s_{4}}^{-1} D_{s_{1} p}, \quad D_{s_{3} s_{1}}^{-1} D_{s_{1}}, 0\right\} \tag{3.9}
\end{equation*}
$$

When conditions $B_{1}\left(\mu_{0}\right)$ or $B_{4}\left(\mu_{0}\right)$ are satisfied, formulas (3.3) do not hold for arbitrary $s_{1}, s_{2}$ and $s_{3}$, system (3.7) has no strictly positive solutions, and (3.6) has no fixed sign integrals, (All integrals have alternating signs when condition $B_{1}\left(\mu_{0}\right)$ is satisfied.)

Let now Rank $C=1$. We assume that $a \neq 0$, and instead of system ( 3,7 ) consider the equation $\quad a \gamma=0$
If condition $B_{3}\left(\mu_{0}\right)$ is satisfied, at least one pair of coefficients $\alpha_{s_{1}}$ and $\alpha_{s_{3}}$ satisfy at point $\mu_{0}$ the equality $\quad \operatorname{sign} a_{s_{1}} a_{s z}=-1$
which is valid in $D$. The general solution of system (3.10) with condition (3,11) is of the form

$$
\begin{equation*}
\gamma_{s}=-a_{s_{1}}^{1}\left(a_{s_{2}} \gamma_{s_{2}}+\sum_{s \neq s_{1}, s_{2}} \gamma_{s} a_{3}\right) \tag{3,12}
\end{equation*}
$$

Selecting $\gamma_{s}\left(s \neq s_{1}\right)$ in conformity with the condition

$$
\gamma_{s}>0 \quad\left(s \neq s_{1}, s_{2}\right), \quad \gamma_{s_{1}}>\max \left\{-\sum_{s \neq s_{1}, s_{2}} \gamma_{s} a_{8}, \mid \rho\right\}
$$

we obtain a strictly positive solution of Eq. (3.10) and a fixed sign integral of system (3.1). When condition $B_{2}\left(\mu_{0}\right)$ is satisfied, condition (3.11) is violated for any arbitrary $s_{1}$ and $s_{2}$, and (3.10) has no positive solution. By selecting all parameters $\gamma_{s}(\mu)$ in the form of continuous functions in $D$ it is possible to assert that the derived integrals are continuous with respect to $\mu$.

Using the set of continuous integrals, we obtain below the criteria of strong stability of system (2.1).
4. Strong atability, Bifurcation. Besides system (2. 10), which is equivalent as regards stability to system (2.1), we consider in region $D$ also the system

$$
\begin{align*}
& u_{s}^{* *}=\lambda_{s}(\mu) u_{s}^{*}+u_{s}^{*} \sum_{|p|=N} \alpha_{s}^{p}(\mu) \omega^{* p}+O_{\mu}^{*}\left(\|u+v\|^{2 N+2}\right)  \tag{4.1}\\
& v_{s}^{* *}=-\lambda_{s}(\mu) v_{s}^{*}+v_{s}^{*} \sum_{|p|=N} \bar{\alpha}_{s}^{p}(\mu) \omega^{* p}+O_{\mu}^{*}\left(\|u+v\|^{N+2}\right)
\end{align*}
$$

which is equivalent to system (2.1) in region $D^{*}$. System (4.1) is obtained by the usual normalization of system ( 2,1 ) in region $D^{*}$.

As shown in Sect. 2, coefficients $\alpha_{s}{ }^{p}$ in systems (2.10) and (4.1) are the same for $k \geqslant 3$. It is important to note that when in (2.10) the nonlinearities $O_{\mu}$ are continuous in $D$, then owing to the coefficients at terms of order not lower than $2 k-1$ in system (4.1) $O_{\mu}{ }^{*} \rightarrow \infty$ when $\mu \rightarrow \mu_{0}$. The coefficients $\alpha_{8}{ }^{p}(|p|=1)$ have this property even at $k=2$. Because of this (4.1) is considered here only for $k>2$.

Let system (3.1) have a continuous set of integrals (3.6). For calculating the derivative of function $V$ we take into account that $V$ is the integral of system (3.1) and, by virtue of system (2.10), we have

$$
\begin{align*}
& \frac{1}{2} V^{*}=\sum_{|q|=N+1} T_{q}(\mu) \omega^{q}+O_{\mu}\left(\|\omega\|^{N+\alpha_{1}}\right)=W_{N+1}(\mu, \omega)+O_{\mu}  \tag{4.2}\\
& T_{q}(\mu)=\sum_{s=1}^{n} \gamma_{s}(\mu) \operatorname{Re} \alpha_{s}^{q-\delta_{s}}
\end{align*}
$$

If the $(N+1)$-st order form $W_{N+1}\left(\mu_{0}, \omega\right)=W_{N+1}^{0}$ is of fixed sign in cone $\omega \geqslant$ 0 , then for a reasonably small $\varepsilon$ function $V^{\prime \prime}$ is of fixed sign in the aeighborhood of zero for all $\mu \in D=L_{z}\left(\mu_{0}\right)$.

We denote by $\Gamma_{0}$ the set of vectors $\gamma\left(\mu_{0}\right)$ for which the form $W_{N+1}^{0}$ is negative definite (obviously $\Gamma_{0}$ can also be $\varnothing$ ). We further denote by $M_{0}$ the set of solutions of system (3.7) when $\mu=\mu_{0}$, and by $M_{0}{ }^{+}$the set of strictly positive solutions. Note that when conditions $A\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$ are satisfied, $M_{0}{ }^{+} \neq \varnothing$.

Theorem 4.1. Let system (2.10) be such that: (1) Rank $C$ is maintained in $D$ for reasonably small $\varepsilon$, and (2) $\Gamma_{0} \neq \varnothing$. Then
a) system (2.10) is strongly asymptotically stable at point $\mu_{0}$, if conditions $A\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$ are satisfied and $\Gamma_{0} \cap M_{0}^{+} \neq \varnothing$;
b) system (2,10) is strongly unstable at point $\mu_{0}$, if one of conditions $\boldsymbol{A}\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$ is satisfied and $\Gamma_{0} \cap\left(M_{0} \backslash M_{0}{ }^{+}\right) \neq \varnothing$.

The validity of this theorem follows from Theorem 3.1 and the theorems of Liapunov's
second method. Let us, for example, consider case (a). Condition (2) implies the existence of vector $\gamma^{\circ}$ such that the form $W_{N+1}^{\circ}$ is negative definite, Condition(a) ensures the existence in set (3.5) of fixed sign integral $V_{0}$ of system(3.1) whose derivative by virtue of $(2,10)$ is negative definite (which is determined by the form $W_{N+1}^{\circ}$ ). Condition (1) and Theorem 3.1 allow us to assert the existence of Liapunov's function that is continuous with respect to $\mu$ and satisfies at every point of region $D$ the conditions of Liapunov's theorem on asymptotic stability.

Case (b) in which the alternating sign integrals of system (3.1) are used, is similarly considered.

Let us now deal with the question of properties of system (2.10) in certain cases when the conditions of Theorem 4.1 are violated. For this we shall consider system (4.1) in $D^{*}$ and the function

$$
V^{*}=\sum_{s=1}^{n} \gamma_{s}(\mu) \omega_{s}^{*}
$$

where $\gamma_{s}$ are arbitrary continuous functions of $\mu\left(V^{*}\right.$, unlike $V$ in (3.6), is not necessarily an integral of system (3.1)).
Taking into account that $\alpha_{s}{ }^{p}$ are the same in systems $(2,10)$ and $(4,1)$, we have

$$
\begin{equation*}
V^{*^{\cdot}}=W_{N+1}\left(\mu, \omega^{*}\right)+O_{\mu^{*}}\left(\|\omega\|^{N+3 / 2}\right) \tag{4.3}
\end{equation*}
$$

where $W_{N+1}\left(\mu, \omega^{*}\right)$ is a form analogous to (4.2) and $O_{\mu}{ }^{*} \rightarrow \infty$ when $\mu \rightarrow \mu_{0}$.
Let the set $\Gamma_{0}$ for form $W_{N+1}\left(\mu, \omega^{*}\right)$ be nonempty. If $\Gamma_{0}$ contains the strictly positive vector $\gamma^{\circ}$, system (4.1) is asymptotically stable in $D^{*}$. Independently of this property of system (4.1), any of the cases $A$ and $B_{j}$ can be realized at point $\mu_{0}$ for system (3.1). If it is assumed that conditions $B_{1}\left(\mu_{0}\right)$ or $B_{2}\left(\mu_{0}\right)$ are satisfied, system (2.10) is then unstable when $\mu=\mu_{0}$ [2]. It is obvious that on these assumptions $\Gamma_{0} \cap M_{0}=$ $\varnothing$, and none of conditions (a) and (b) are satisfied.

The above considerations lead to the following result.
Theorem 4.2. If system (2.10) is such that conditions $B_{1}\left(\mu_{0}\right)$ or $B_{2}\left(\mu_{0}\right)$ are satisfied, and form $W_{N+1}\left(\mu_{0}, \omega^{*}\right)$ can be made negative definite, then $\mu_{0}$ is a bifurcation point at which the asymptotic stability in $D^{*}$ changes to instability. If the set $\Gamma_{0} \neq \varnothing$ does not contain strictly positive vectors, then system (2.10) is strongly unstable in $D$ when conditions $B_{1}\left(\mu_{0}\right)$ or $B_{2}\left(\mu_{0}\right)$ are satisfied.

Note that the region in which function $V^{* *}$ is of fixed sign contracts under conditions of Theorem 4,2 when $\mu \rightarrow \mu_{0}$ to the coordinate origin, as implied by the properties of function $O_{\mu^{*}}$ in (4.3). When $\mu=\mu_{0}$ an instability region, lying in the neighborhood of the unstable solution of system (3.1), is generated [2].

Example. Let us consider the system of differential equations

$$
\begin{align*}
& z_{s} \because+\rho_{s}^{2}(\mu) z_{s}=Z_{s}^{(2 N)}\left(\mu, z, z^{\cdot 2}\right)+Z_{s}^{2 N+1}\left(\mu, z, z^{\cdot 2}\right)+g_{s}(\mu) z_{s}{ }^{2 N+1}+\ldots  \tag{4.4}\\
& (s=1,2, \ldots, n, N>1)
\end{align*}
$$

where $\rho_{s}(\mu), g_{s}(\mu)$ and the coefficients of forms $Z_{s}{ }^{(j)}$ are continuous functions of $\mu$.
Using the substitution $x_{s}=z_{s}-\left(i / \rho_{s}\right) z_{s}{ }^{*}=\eta_{s}$ followed by the transformation to the continuous normal form, we reduce the problem to the analysis of system (2.10) whose coefficients are defined as follows:

$$
\lambda_{s}=i_{s}, \quad \alpha_{s}(\mu)=-\frac{i}{\rho_{s}(\mu)} b_{s}(\mu), \quad \alpha_{s}^{p}(\mu)=-\frac{i}{\rho_{s}(\mu)} b_{s}^{p}(\mu) \text { for } p \neq N \delta_{z}
$$

$$
a_{s}^{N \delta_{s}}(\mu)=\rho_{s}^{2 N} 2^{-2 N-1} C_{2 N+1}^{N} g_{s}-\frac{i}{\rho_{s}} b_{s}^{N \delta_{z}}
$$

The real numbers $b_{s}(\mu)$ and $b_{s}^{p}(\mu)$ are coefficients at terms of inner and identicalresonances in forms $Z^{(j)}\left(\mu, 1 / 2\left(x_{s}+y_{s}\right),-1 / 4 \rho_{s}^{2}\left(y_{s}-x_{s}\right)^{2}\right)$ for $j=2 N$ and $2 N+1$, respectively. At point $\mu_{0}$ the numbers $b_{s}\left(\mu_{0}\right)=b_{s}{ }^{\circ}$ and $g_{s}\left(\mu_{0}\right)=g_{s}{ }^{c}$ are nonzero.

It is seen that for system (4.4) matrix $C(\mu)$ retains its rank in $D$, and that at point $\mu_{0}$ either conditions $B_{2}\left(\mu_{0}\right)$ or $B_{3}\left(\mu_{0}\right)$ are satisfied when either all $b_{s}{ }^{\circ} / \rho_{s}{ }^{\circ}$ are of the same sign or when among $b_{8} \% \rho_{s}^{\circ}$ there are numbers of different signs, respectively.

Form $W_{N+1}^{\circ}$ and the sets $\Gamma_{0}$ and $M_{0}$ are defined as follows:

$$
\begin{aligned}
& W_{N+1}^{\circ}=\sum_{s=1}^{n} \gamma_{s}^{\circ} \operatorname{Re} \alpha_{s}^{N \delta_{s}}\left(\mu_{0}\right) \omega_{s}^{N+1} \quad\left(\operatorname{sign} \operatorname{Re} \alpha^{N \delta_{s}}\left(\mu_{0}\right)=\operatorname{sign} g_{\mathrm{s}}{ }^{\circ}\right) \\
& \Gamma_{0}=\left\{\gamma^{\circ} \mid \operatorname{sign} \gamma_{s}^{\circ}=-\operatorname{sign} g_{\mathrm{s}}^{\circ}\right\}, \quad V_{0}=\left\{\gamma^{\circ} \mid \sum_{j=1}^{n} b_{j}^{\circ} \gamma_{j}^{\circ}=0\right\}
\end{aligned}
$$

It is evident that $\Gamma_{0} \neq \varnothing$ and conditions (1) and (2) of Theorem 4.1 are satisfied for system (4.4).

When all $g_{s}^{\circ}<0$ and condition $B_{3}\left(\mu_{0}\right)$ is satisfied, system (4.4) is strongly asymptotically stable (conditions (a) of the theorem are satisfied).

If among $g_{s}{ }^{\circ}$ there are numbers of different signs or all $g_{s}{ }^{\circ}>0$, then for certain $k$, $j$ will be sign $g_{k}{ }^{\circ} b_{k}{ }^{\circ}=-\operatorname{sign} g_{j}{ }^{\circ} b_{j}{ }^{\circ}$, condition (b) of Theorem 4.1 is satisfied, and system (4.4) is strongly unstable. The last requirement is automatically realized when all $b_{s} / \rho_{s}{ }^{\circ}$ are of the same sign and among $g_{s}{ }^{\circ}$ there are numbers of different signs, or vice versa.

If all $g_{s}^{\circ}<0$ and condition $B_{2}\left(\mu_{0}\right)$ is satisfied, then by Theorem 4.2 at point $\mu_{0}$ the bifurcation - the asymptotic stability in $D^{*}$ - changes to instabilit\} at that point.

Finally, if condition $B_{2}\left(\mu_{0}\right)$ is satisfied and $(\forall s) g_{s}{ }^{\circ} b_{s} / \rho_{s}{ }^{\circ}<0(>0)$, by Theorem 4.2 we have strong instability at point $\mu_{0}$.

We note in conclusion that Theorem 4.2 may be amplified as follows.
Theorem 4.3. If conditions $B_{1}\left(\mu_{0}\right)$ or $B_{2}\left(\mu_{0}\right)$ are satisfied for system (2.10) and system (4.1) is asymptotically stable in $D^{*}$ within terms of order not higher than $2 k-2$, then $\mu_{0}$ is the bifurcation point of (2.10) of the same kind as in Theorem 4.2.

If system $(4,1)$ is unstable in $D^{*}$ to within terms of order not higher than $2 k-2$ and conditions $B_{1}\left(\mu_{0}\right)$ or $B_{2}\left(\mu_{0}\right)$ are satisfied, then system (2.10) is strongly unstable at point $\mu_{0}$.

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